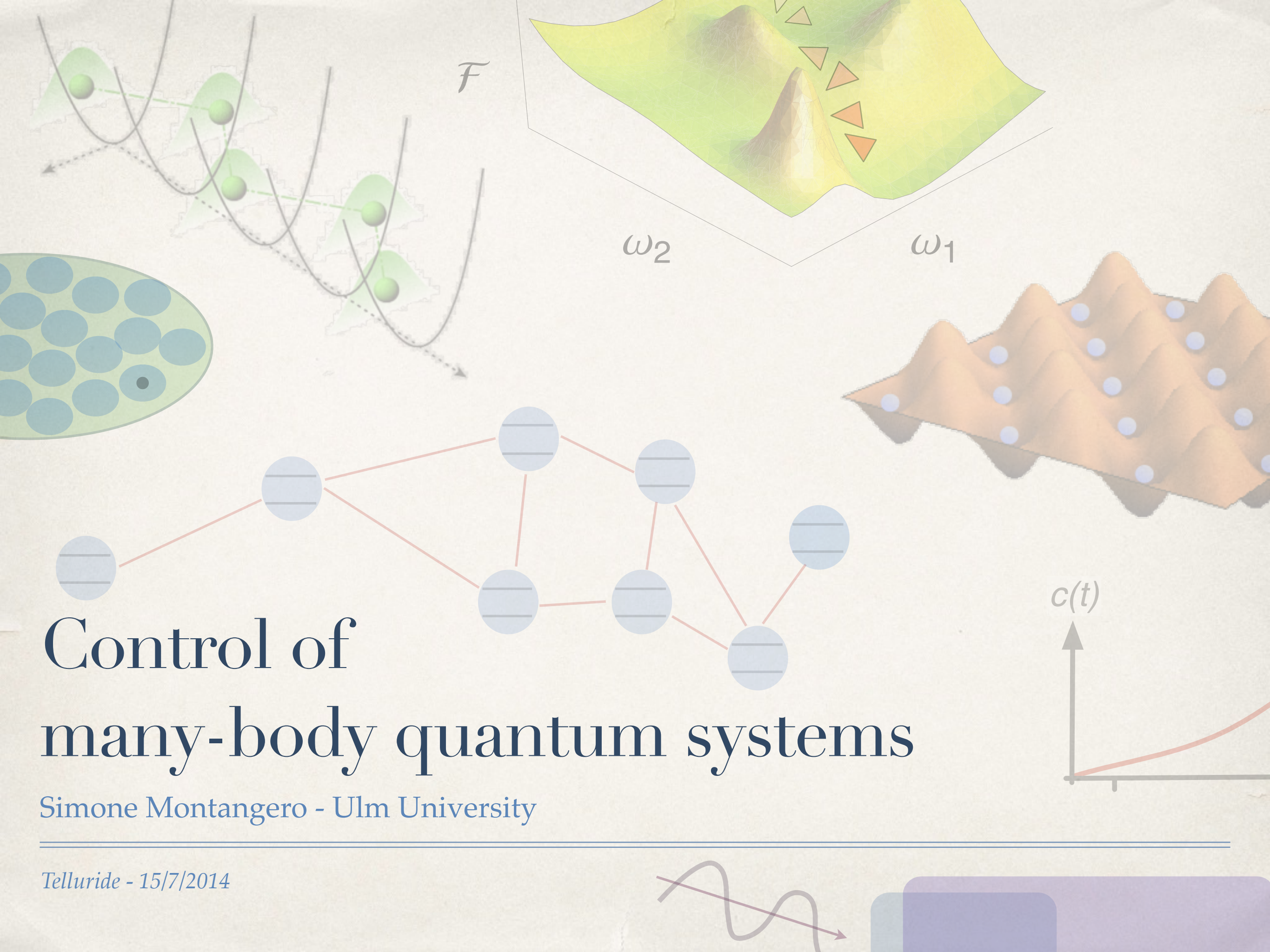




Control of many-body quantum systems

Simone Montangero - Ulm University

Telluride - 15/7/2014

Simulation and control

- Quantum Information Processing (computation, communication etc.)
- Quantum simulators benchmarking and validation
- Design and validation of experiments with quantum hardware
- HEP investigations?



Many body quantum systems

Time dependent Schrödinger equation

$|e\rangle$
 $|g\rangle$

$\mathcal{H}_{loc}$



$\mathcal{H}_{loc}$



$\mathcal{H}_{loc}$

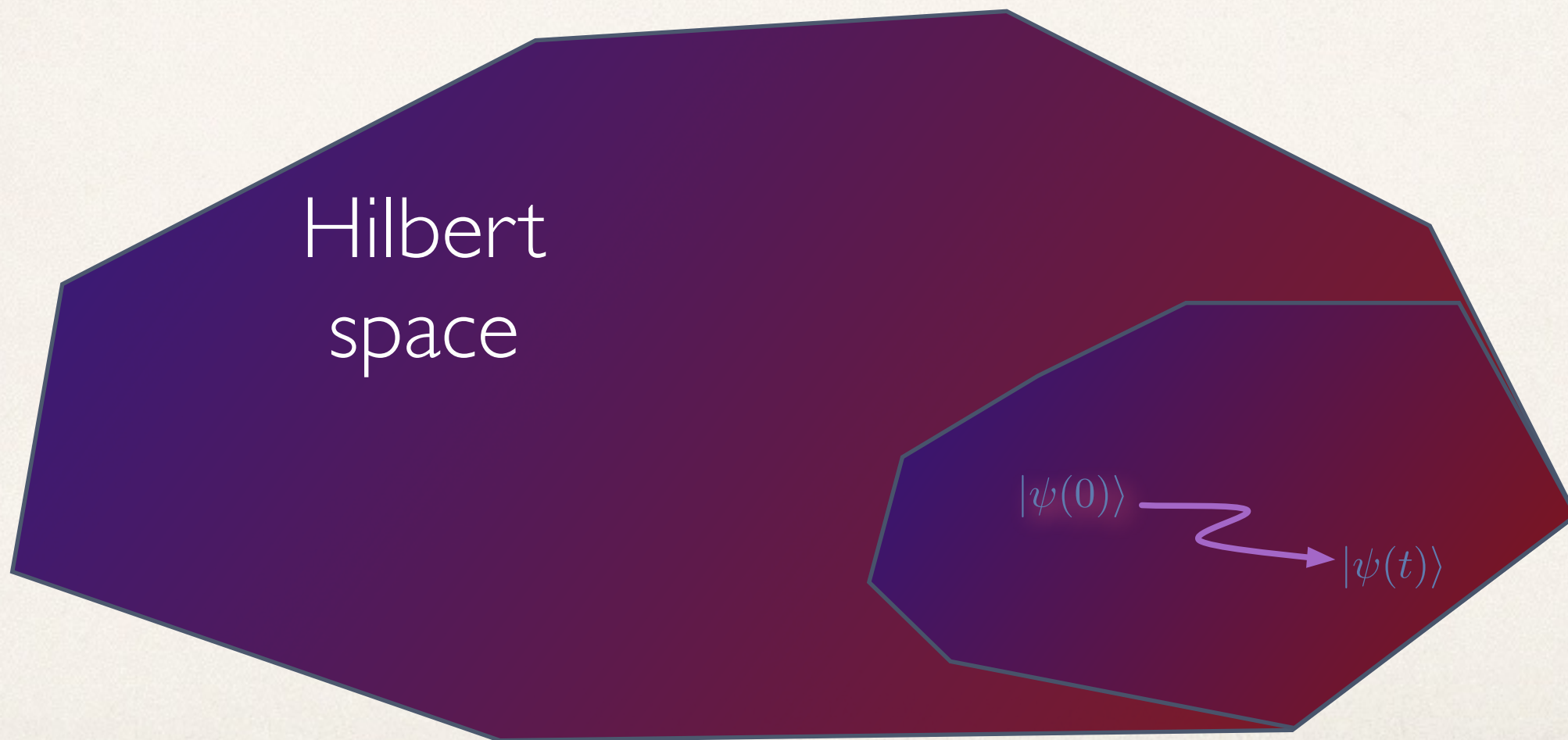
$+ \mathcal{H}_{int}$

$+ \mathcal{H}_{int}$



Tensor Network methods

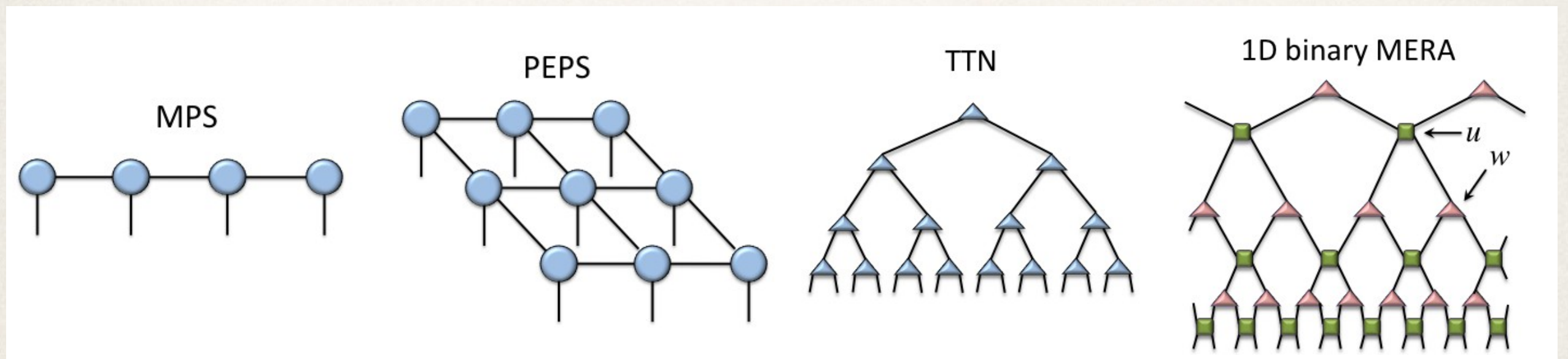
“Simple” Ansatz to describe faithfully “interesting” quantum states!



Tensor Network Methods

$$|\psi\rangle = \sum_{\vec{\alpha}} \psi_{\alpha_1, \alpha_2, \dots, \alpha_n} |\alpha_1\rangle |\alpha_2\rangle \dots |\alpha_n\rangle$$

$|\alpha_i\rangle$
local basis



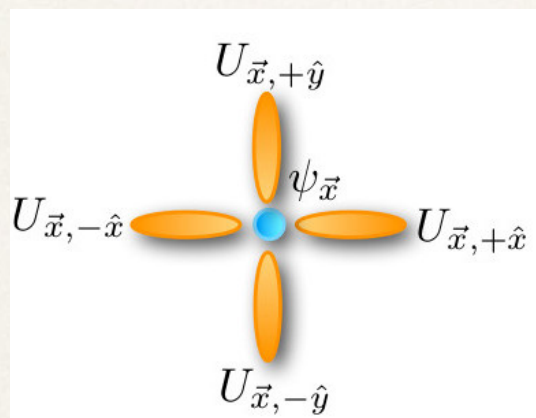
$$|\psi\rangle = \sum_{\vec{\alpha}} A_{\alpha_1}^{\beta_1} A_{\alpha_2}^{\beta_1 \beta_2} \dots A_{\alpha_n}^{\beta_{n-1}} |\alpha_1\rangle |\alpha_2\rangle \dots |\alpha_n\rangle$$

Polynomial efforts
in 1D

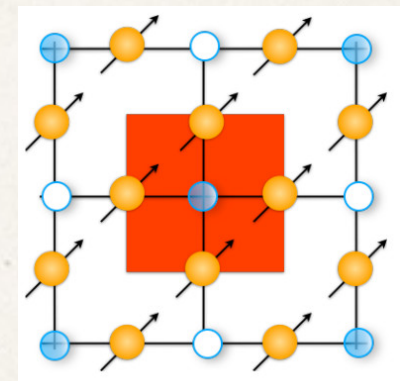
Lattice Gauge Theories

- Site d.o.f. (matter fields..)
- Quantum link d.o.f. (bosons, spins,..)
- Gauge invariant Hamiltonian

Strongly correlated
models at low energies
(spin ice, quantum dimer..)



Non-perturbative
HEP models
(QED, QCD)



Atomic quantum
simulators

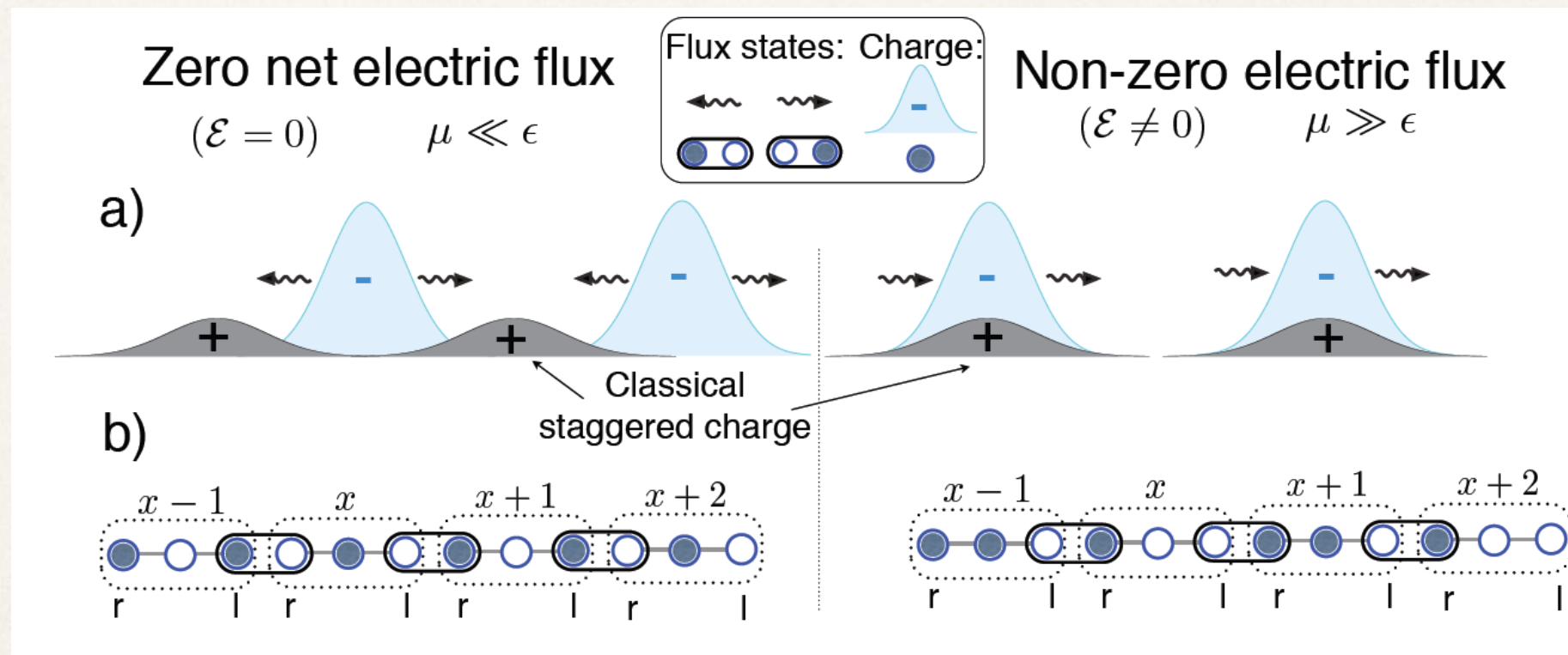
Lattice gauge tensor network
simulations
(abelian and non abelian)

E.Rico, T. Pichler, M. Dalmonte,
P. Zoller, SM, PRL (2014)

P. Silvi, E.Rico, T. Calarco, SM, arXiv:1404.7439

Schwinger model in 1+1D

$$H = \sum_x (E_{x,x+1} - (-1)^x E_0)^2 + \mu(-1)^x \psi_x^\dagger \psi_x - \epsilon(\psi_x^\dagger U_{x,x+1} \psi_{x+1} + \text{h.c.})$$



Disordered phase

Quantum phase
transition

Ordered phase
symmetry breaking

Optimal control



Few-body quantum systems: standard optimal control
(high-accuracy, complete knowledge, many iterations...)

* Many-body?

H. Rabitz, NJP (2009)
Altafini & Ticozzi IEEE (2012)

Chopped RAndom Basis

Expand control field over n_f
“randomized” basis functions

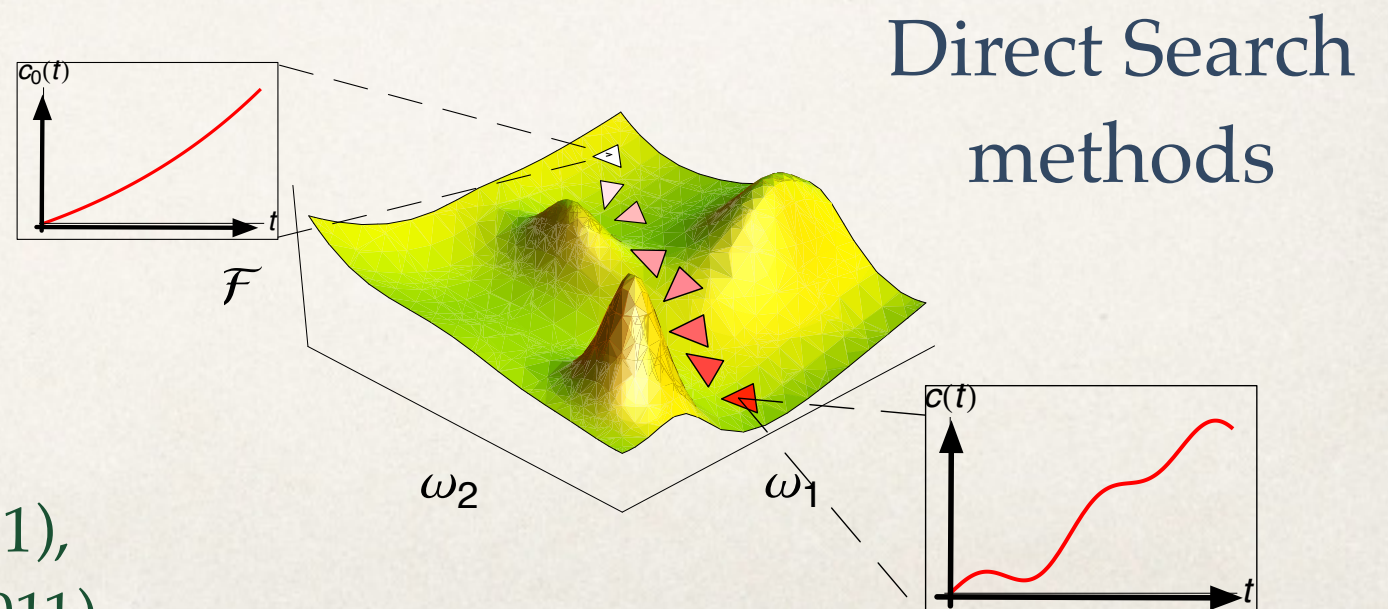
Reduced basis method

Multivariable
function minimization

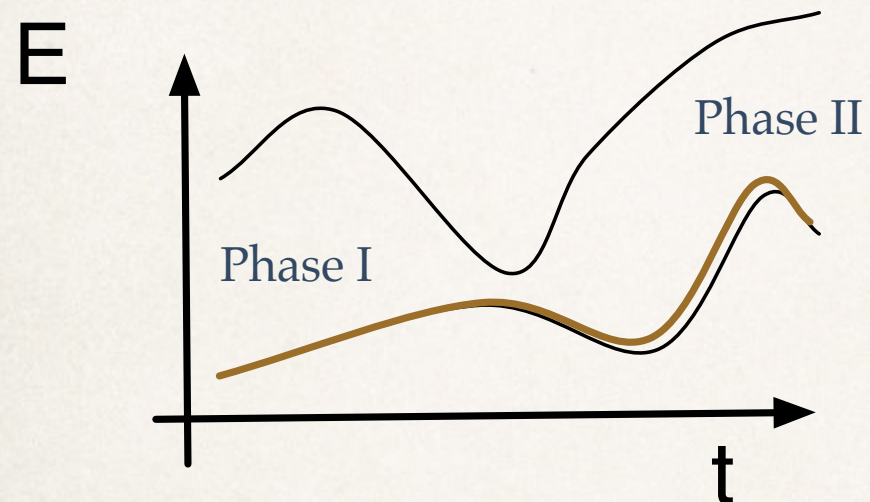
Functional
minimization

$O(10)$ parameters!

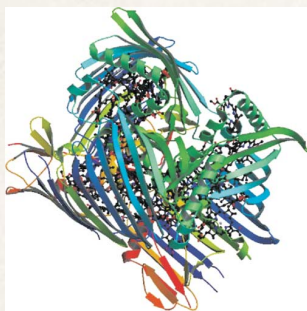
P. Doria, T. Calarco, SM PRL (2011),
T. Caneva, T. Calarco, SM, PRA (2011)



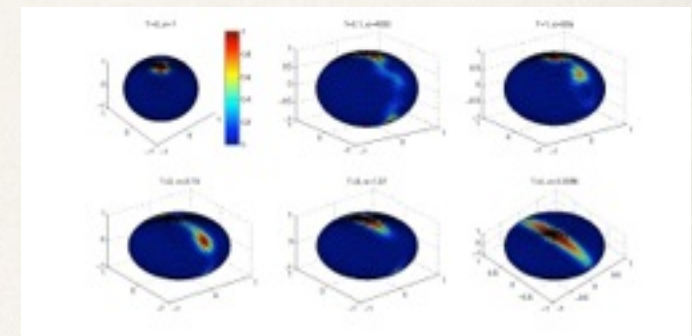
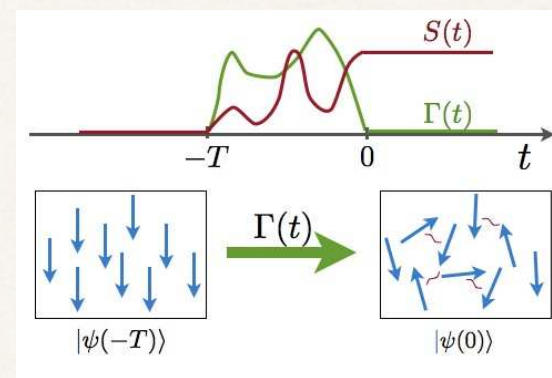
Applications



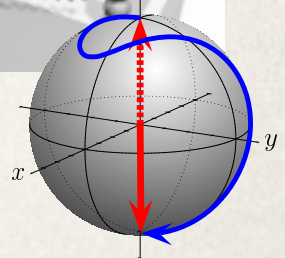
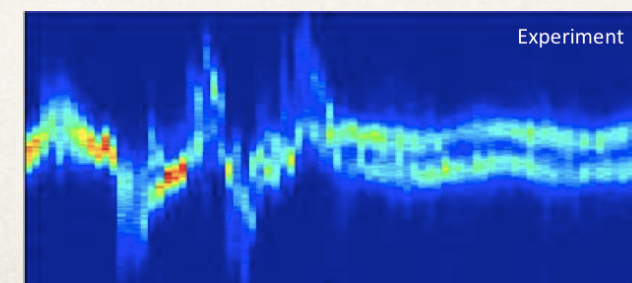
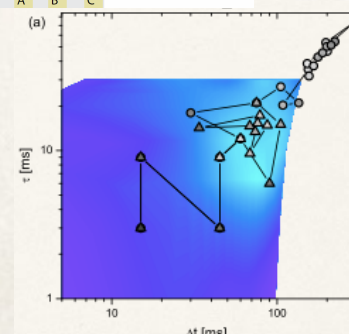
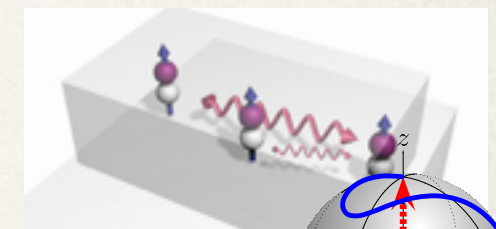
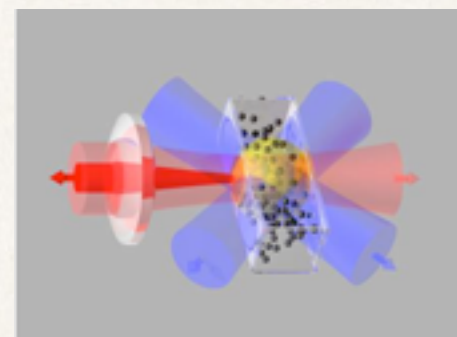
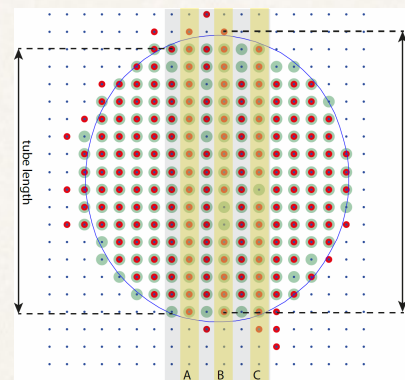
Quantum Phase Transition
dynamics



Light-harvesting
dynamics



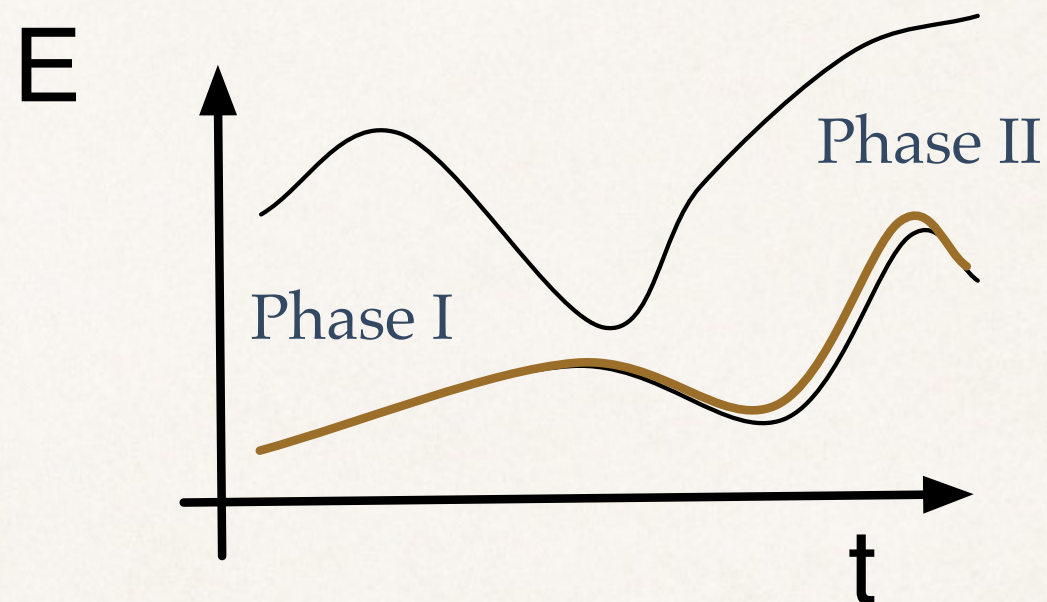
Entanglement/Squeezing manipulation



Optimal experimental protocols

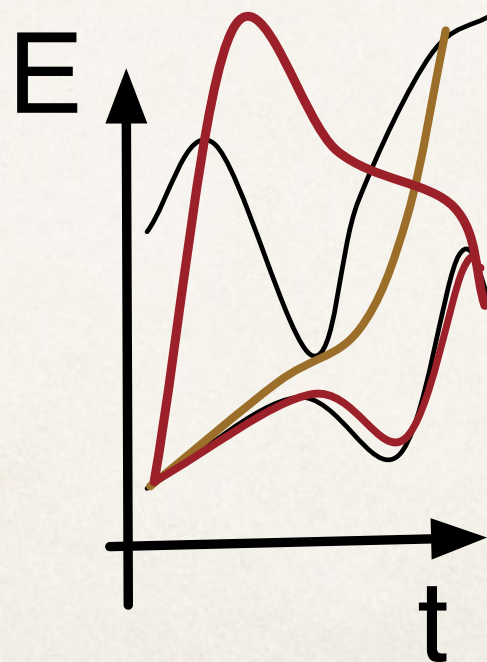
Control of QPT dynamics

Slow



Adiabatic
strategy

Fast

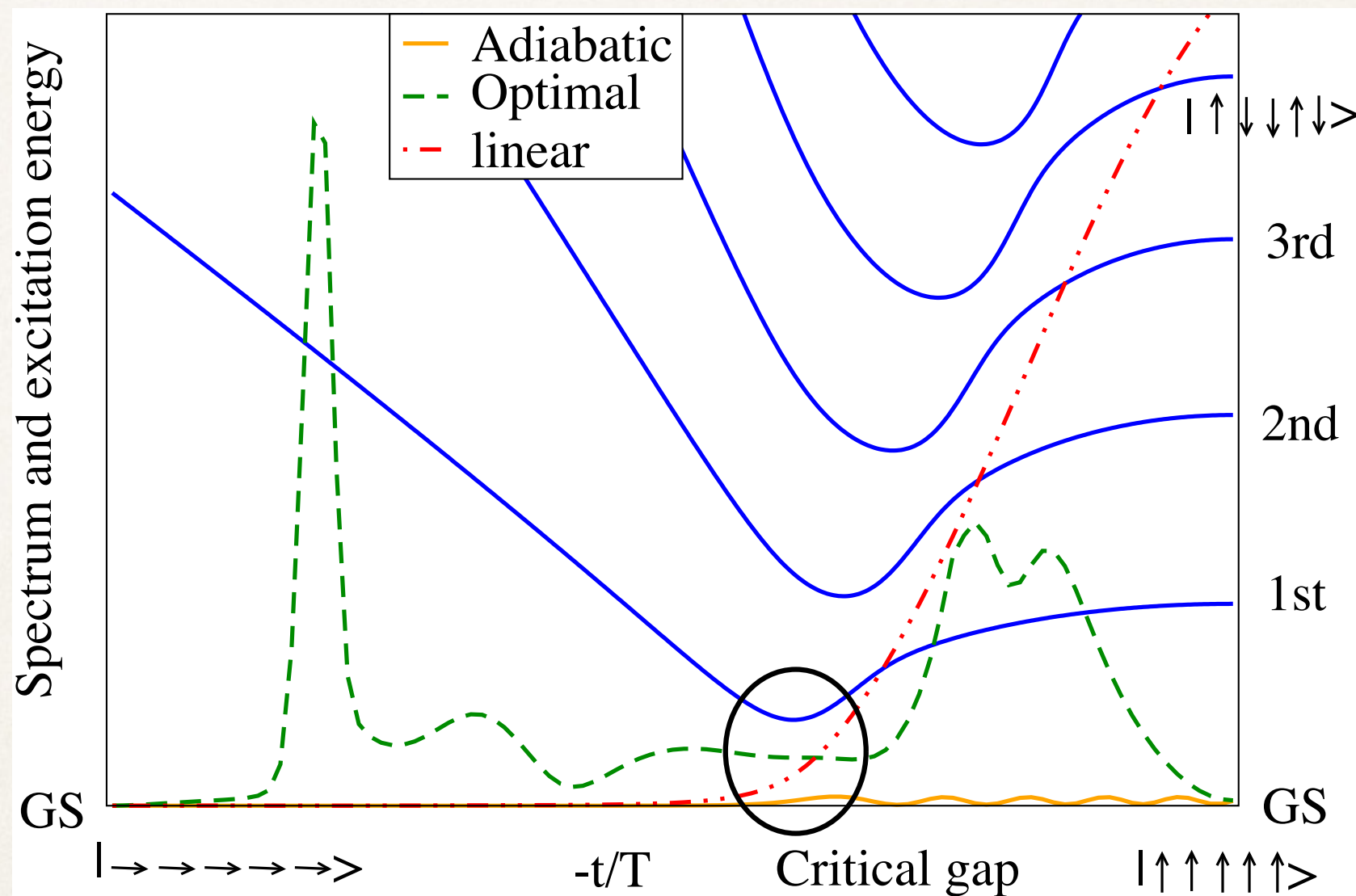


Shortcuts

Optimal
control

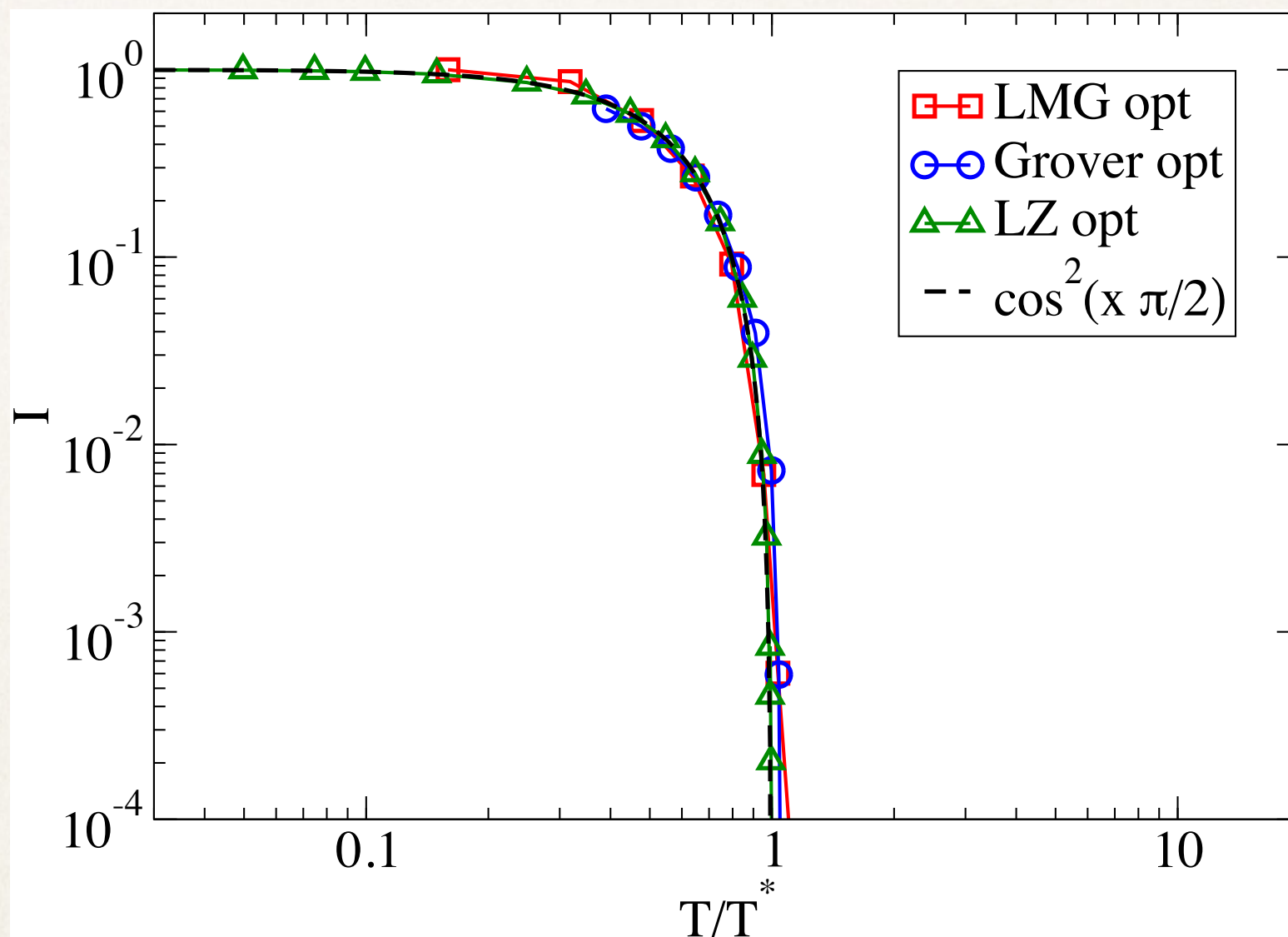
Optimal QPT crossing

LMG
model



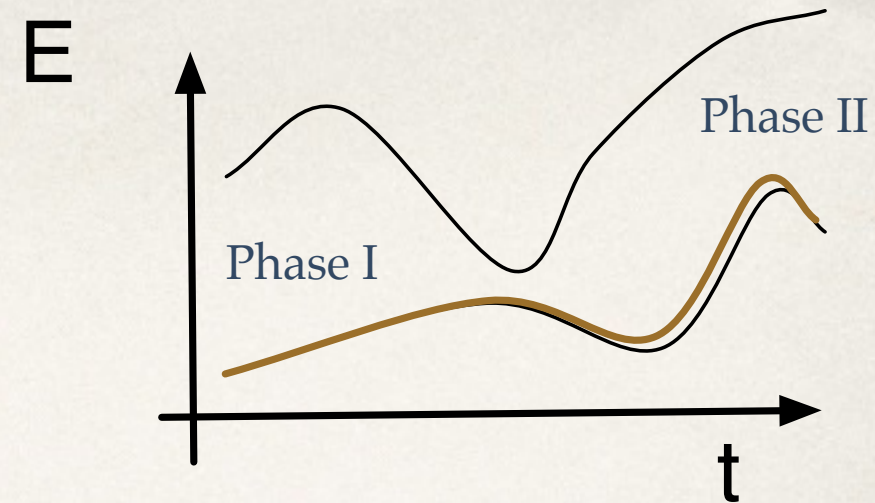
$$H = -\frac{1}{N} \sum_{i < j} (\sigma_i^x \sigma_j^x + \gamma \sigma_i^y \sigma_j^y) - \Gamma(t) \sum_i \sigma_i^z$$

Quantum speed limit



see also T. Caneva, M. Murphy, T. Calarco, R. Fazio, SM, V. Giovannetti, and G. E. Santoro, Phys. Rev. Lett. 103, 240501 (2009).

QPT crossing scenario



$$T \ll \Delta^{-1}$$

$$T \sim \Delta^{-1}$$

$$T \gg \Delta^{-1}$$

Linear

s

$$I \sim O(1)$$

$$I(s \gg 1) \rightarrow 0$$

π

∞

Optimal

s

$$I = \cos^2(s/2)$$

$$I = 0$$

π

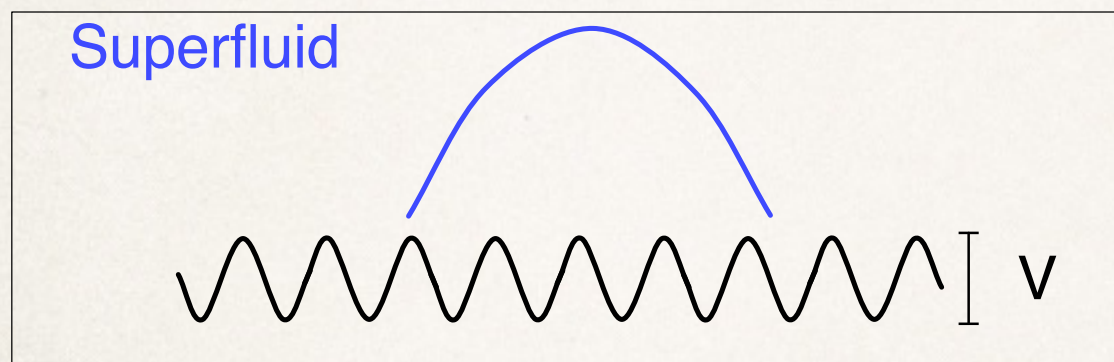
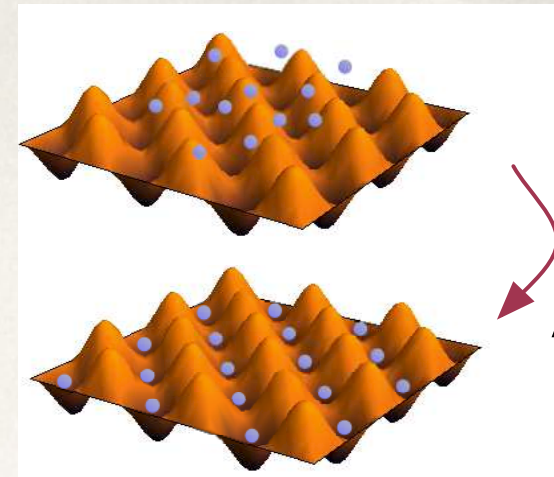
∞

Kibble Zurek

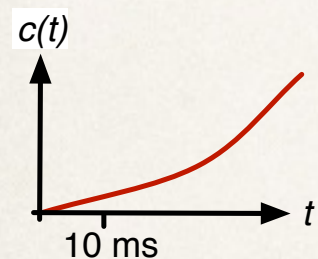
QSL

Adiabatic

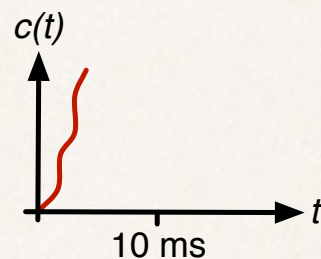
Open loop optimization



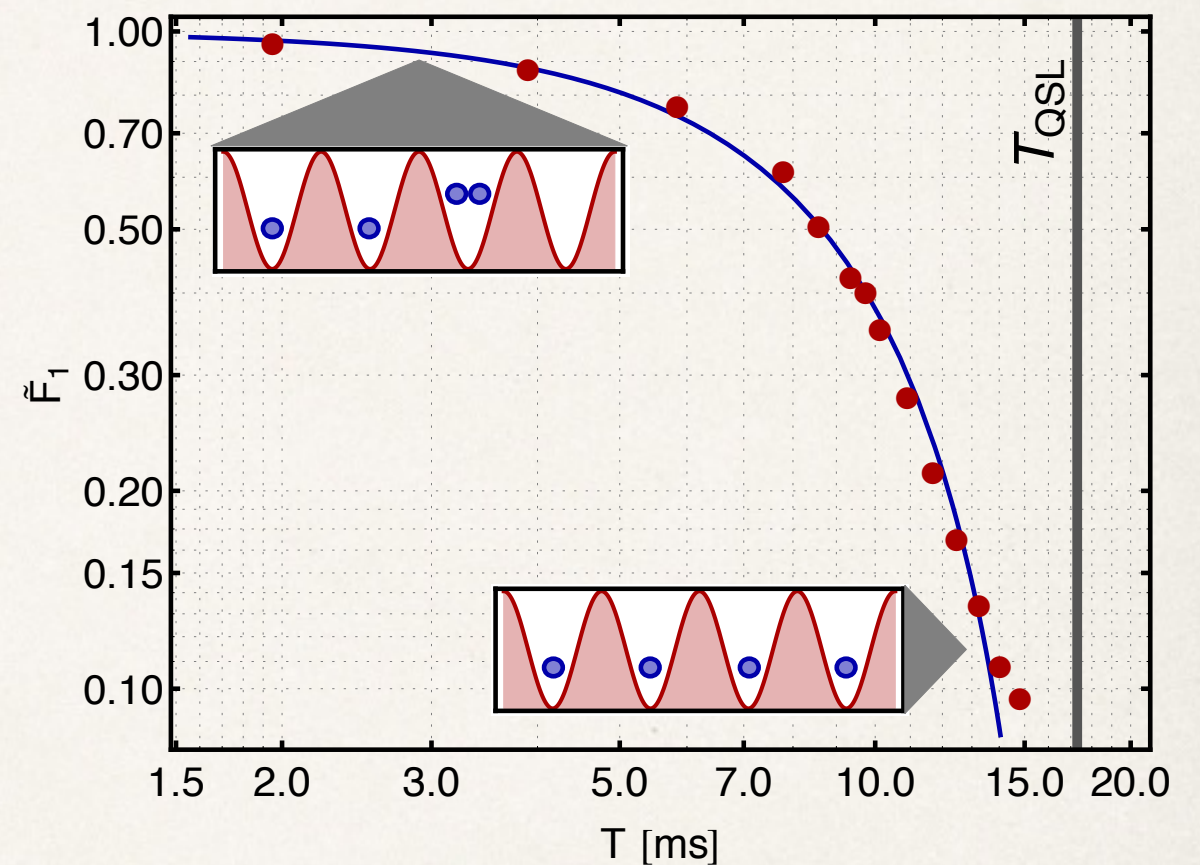
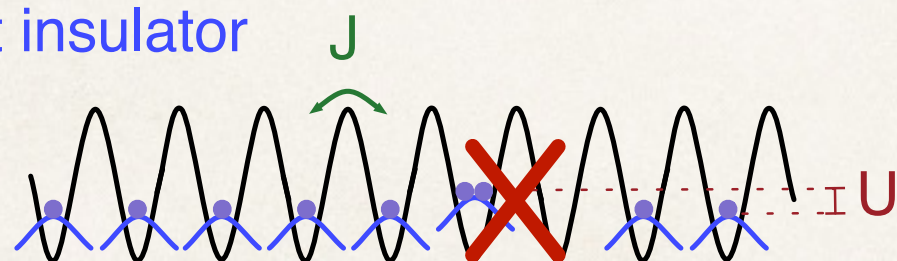
Adiabatic



Antiadiabatic



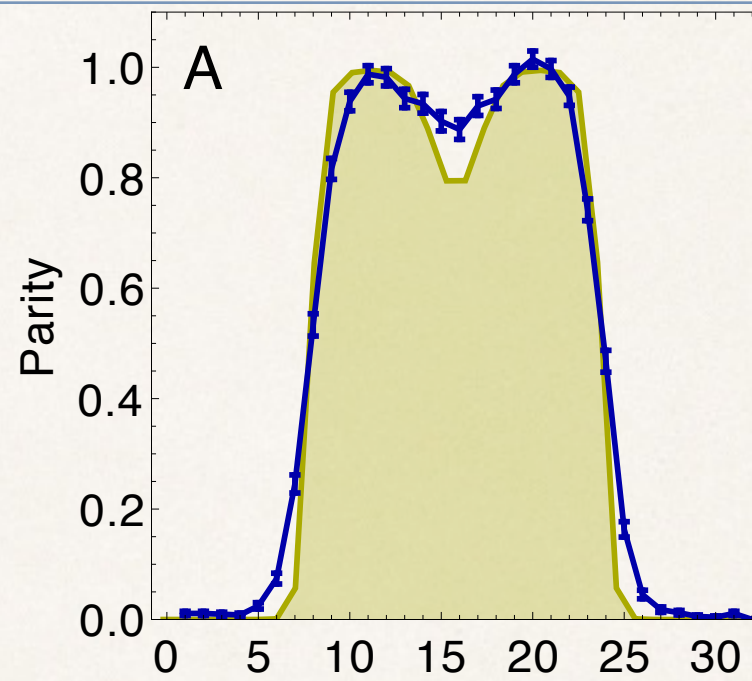
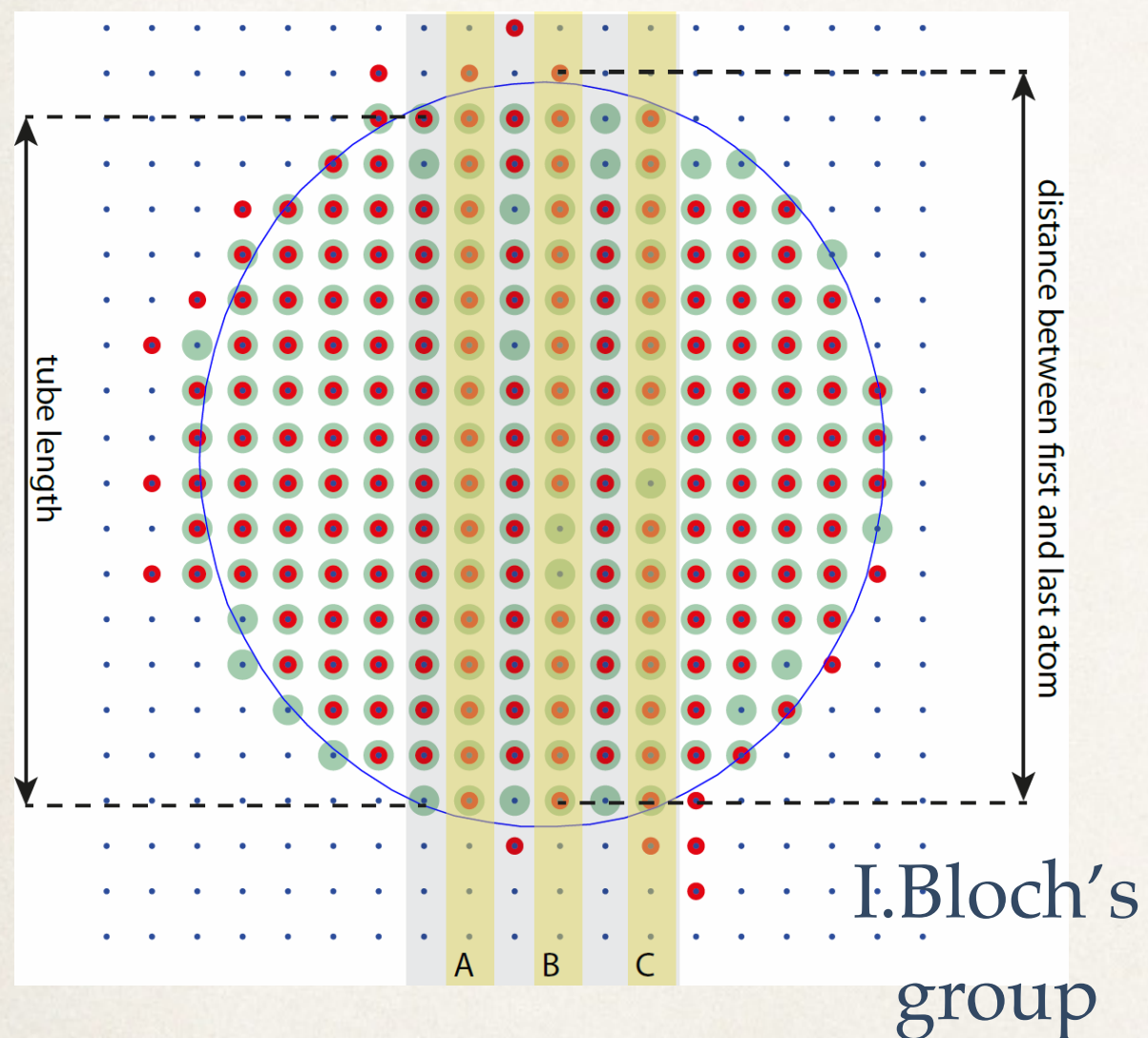
Mott insulator



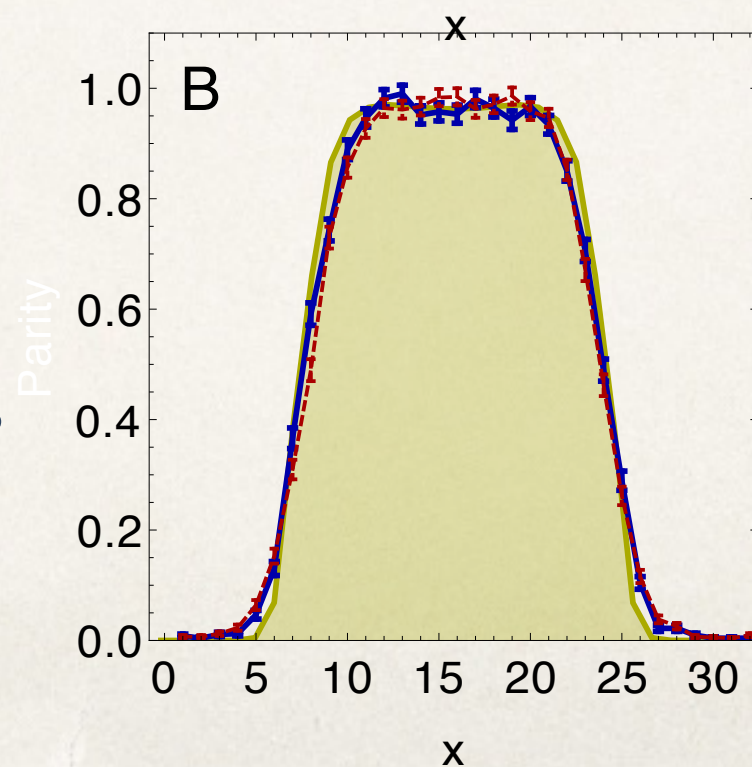
Quantum speed limit

$$T_{QSL} \sim 11 - 15 \text{ ms}$$

QPT crossing at QSL



Linear ramp
 $T=12\text{ms}$



Optimal ramp
 $T=12\text{ms}$
Adiabatic ramp
 $T=110\text{ms}$

Optimal control limits

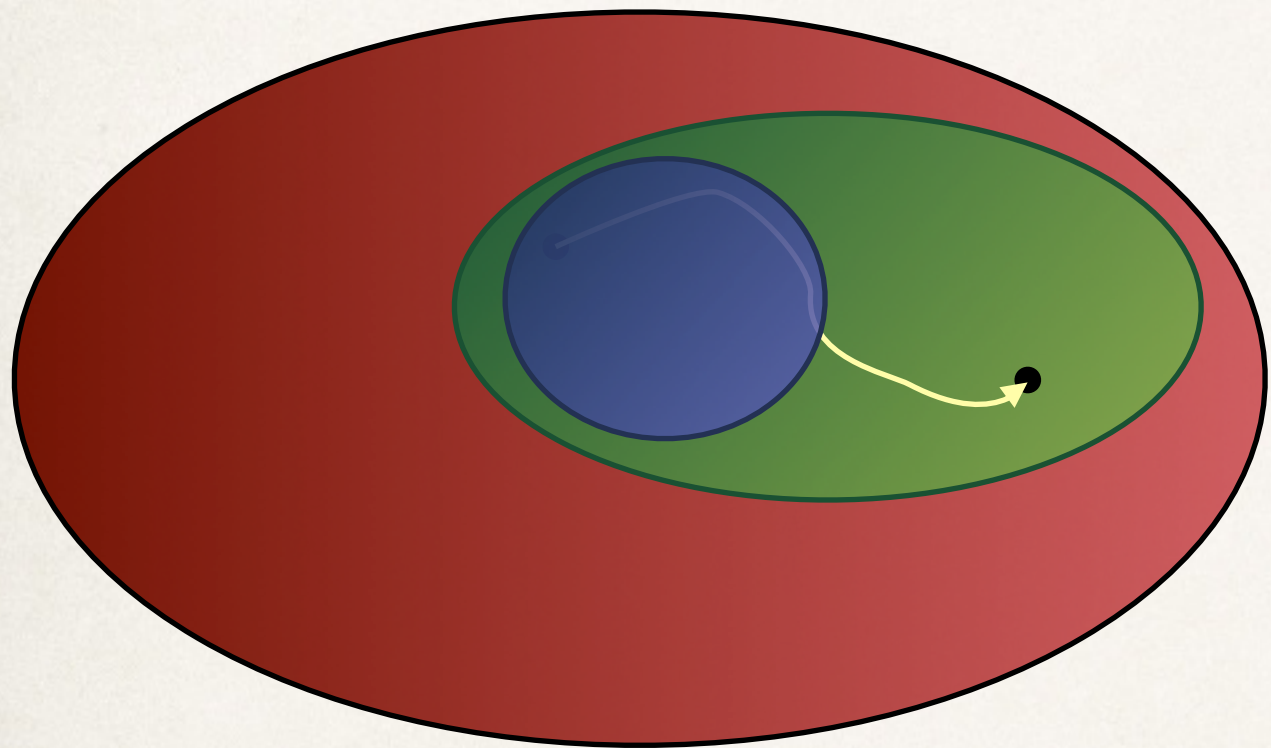
- ❖ What are the physical limits of control of MBQS?

Controllability, Reachability, Quantum Speed Limit...

Control complexity

- ✧ What are the physical limits of control of MBQS?
Controllability, Reachability, Quantum Speed Limit, ...
- ✧ Are there any algorithmic/informational limits?
- ✧ How to characterize the complexity of the optimization task?

Reachable states



$$\mathcal{H} = \mathbb{C}^D$$

Reachable states $\mathcal{W}$

$$\|H_D\| = \|H_C\| = 1$$

$$\gamma(t) \in [0 : 1] \forall t$$

$$H = H_D + \gamma(t)H_C$$

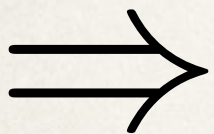
$$\dot{\rho} = \mathcal{L}(\rho, \gamma(t))$$

Time-poly Reachable states $\mathcal{W}^+$

$$D_{\mathcal{W}^+}(N)$$

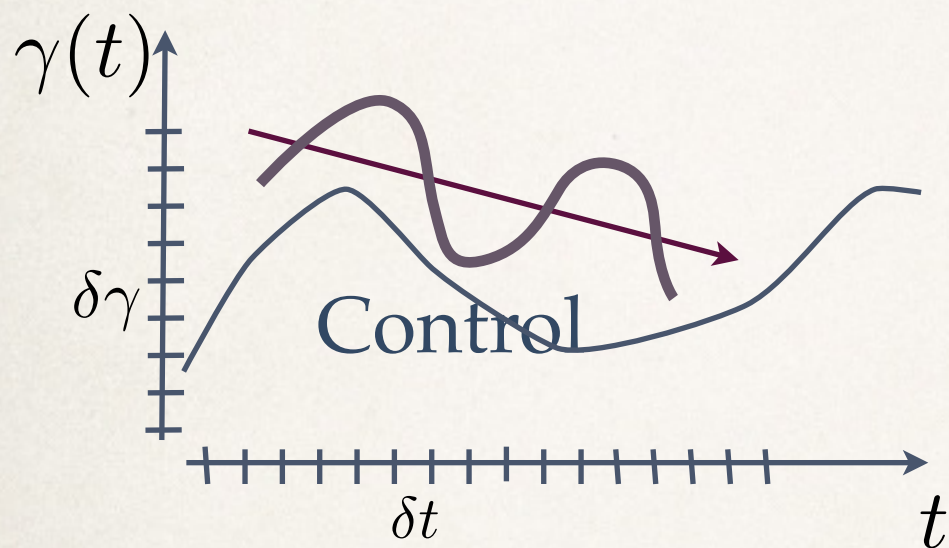
Optimal control complexity

The (smoothed) complexity of an optimal control problem
is **polynomial** in the dimension
of the set of Time-poly reachable states $D_{\mathcal{W}^+}(N)$



The optimal control problem can be solved efficiently via CRAB

Lower Bound

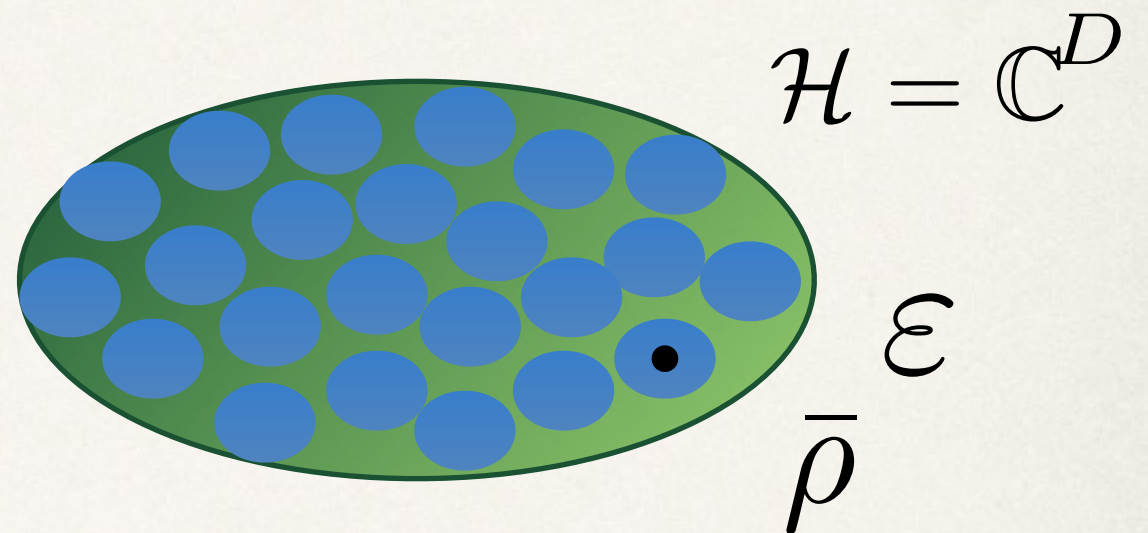


Carried bits:

$$b_\gamma = T \Delta\Omega \kappa_S$$

$$b_\gamma \geq b_S^-$$

$$\varepsilon \geq 2^{-\frac{\kappa n_S}{D}}$$



Minimal needed bits:

$$b_S^- = \log \varepsilon^{-D}$$

$$n_S \geq D$$

$$n_S = \text{Poly}(D_{\mathcal{W}+})$$

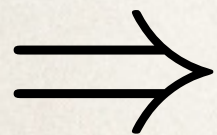
Similar arguments lead to a polynomial upper bound...

CRAAB smoothed complexity

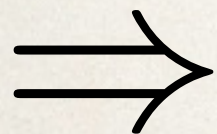
The simplex algorithm minimize the cost function

$$F(n_s) = F(\text{Poly}(D_{\mathcal{W}+}(N)))$$

its smoothed complexity is polynomial in $D_{\mathcal{W}+}$



Integrable and TN-simulatable dynamics can be efficiently controllable.



General non-Integrable systems are exponentially difficult to control.

Optimal control complexity

- ❖ t-DMRG allows to simulate dynamics of many-body quantum systems that can be described efficiently by a MPS state.
- ❖ A time-dependent MPS lives in a space which is $\text{Poly}(n)$, thus the optimal control problem dimension is at most polynomial in n .
- ❖ Slightly entangled dynamics can be efficiently represented via MPS, thus can be also controlled.
- ❖ In general, what can be simulated / represented efficiently can also be optimally controlled!

Time bounds

Lower bound:

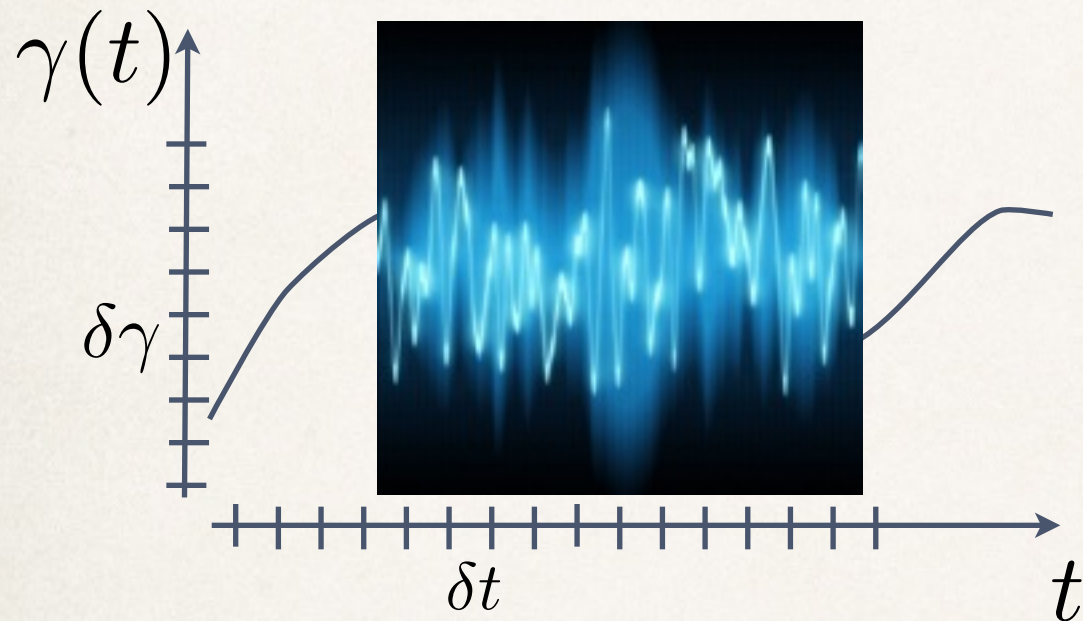
$$\varepsilon \geq 2^{-\frac{T \Delta\Omega \kappa_S}{D_{\mathcal{W}}}} \implies T \geq \frac{D_{\mathcal{W}}}{\Delta\Omega \kappa_S} \log(1/\varepsilon)$$

$$\Downarrow \kappa_{\varepsilon} = \kappa_S$$

Continuous-time
Solovay-Kitaev theorem

$$T \geq \frac{D_{\mathcal{W}}}{\Delta\Omega}$$

Noise effects

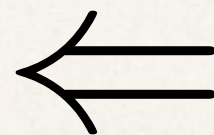


~~$$\kappa_S = \log(1 + \gamma_{max}/\delta\gamma)$$

bits per unit time~~

$$\varepsilon \geq (1 + S/N)^{-\frac{n_s}{D_{\mathcal{W}}}}$$

$$\varepsilon \gtrsim (N/S)^{n_s/D_{\mathcal{W}}}$$

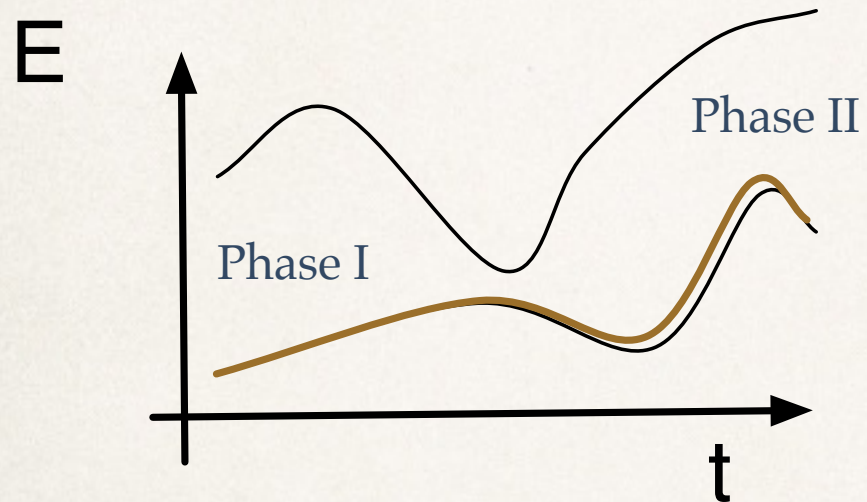


Shannon-Hartley theorem

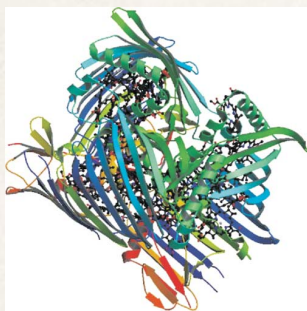
$$k_S = \log(1 + S/N)$$

If $n_s = D_{\mathcal{W}}$ linearly sensitive to noise!

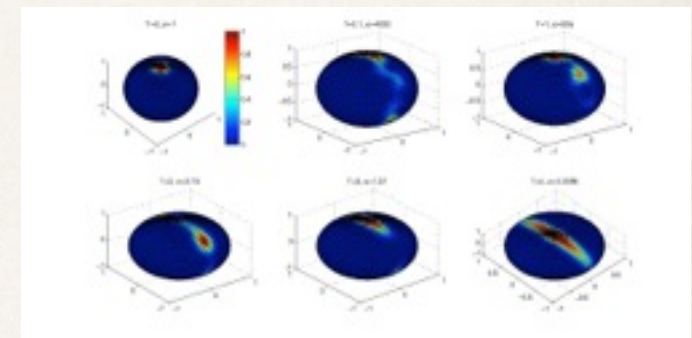
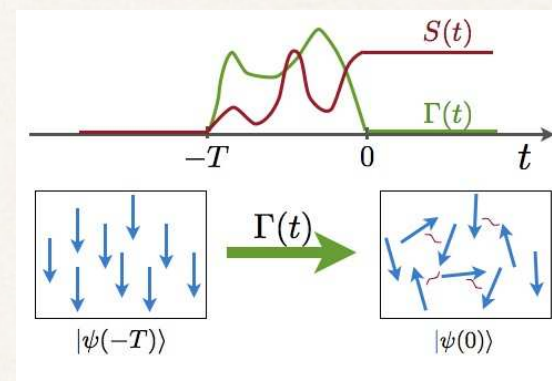
Applications (continued)



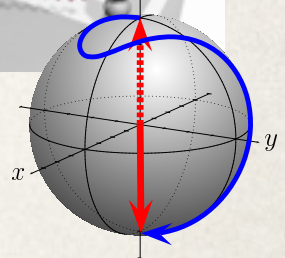
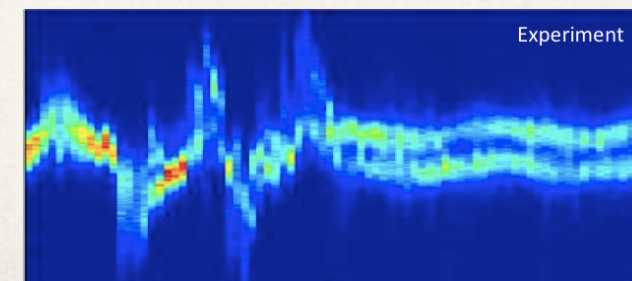
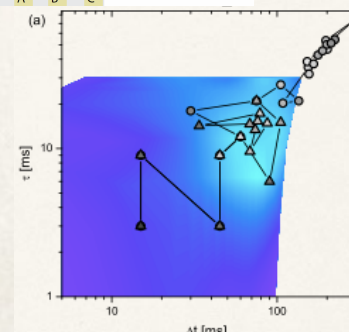
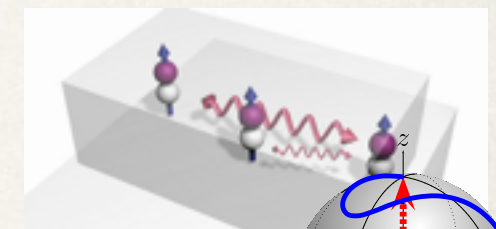
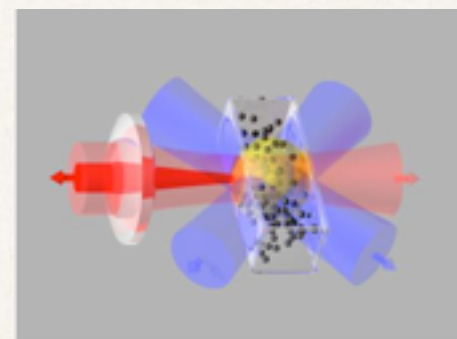
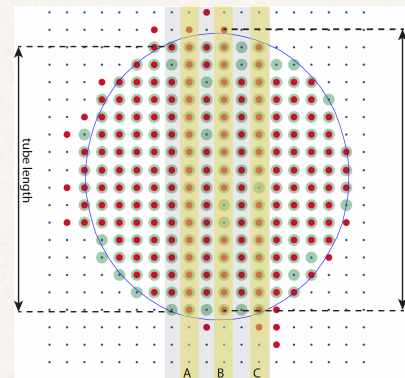
Quantum Phase Transition
dynamics



Light-harvesting
dynamics

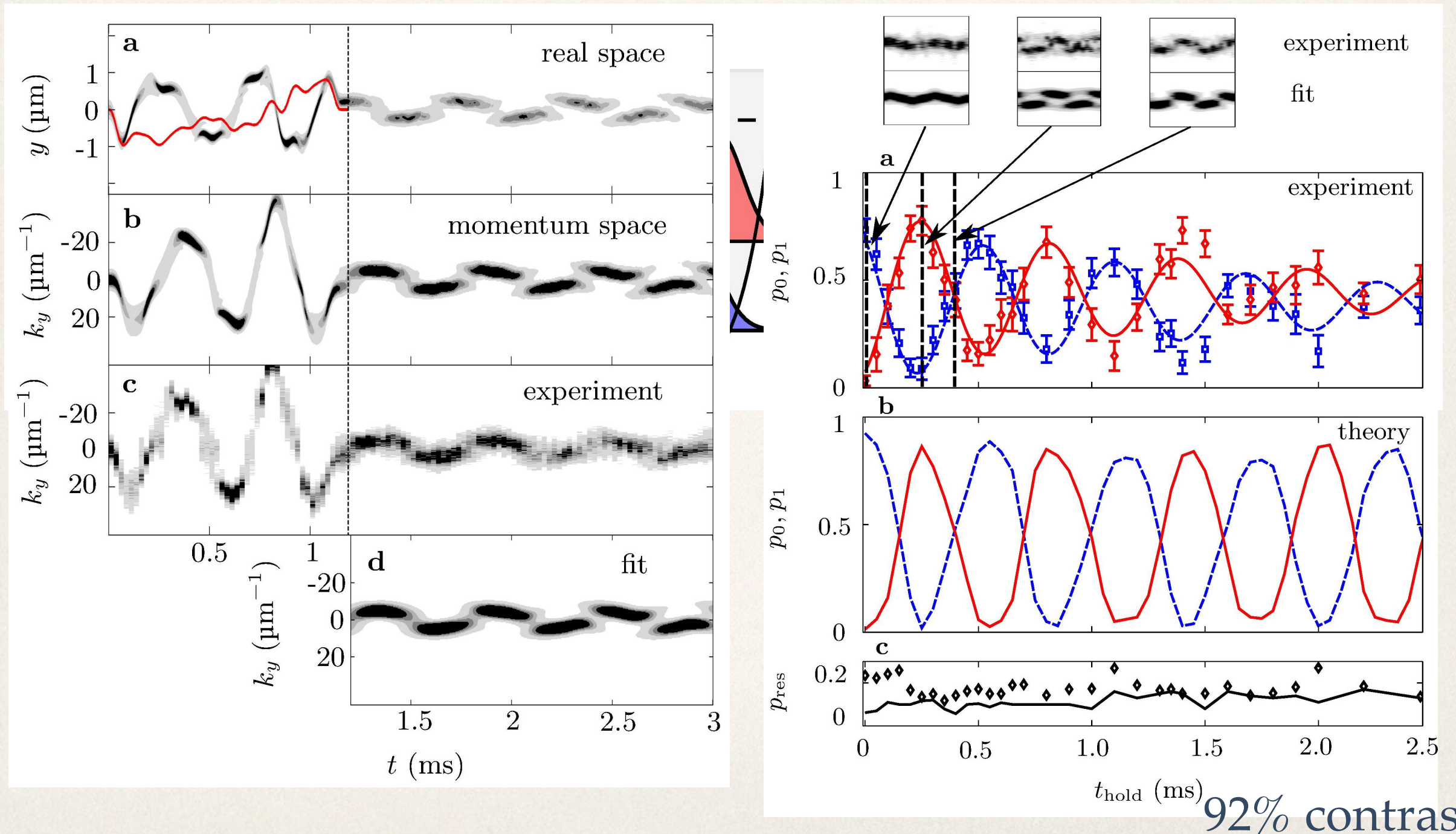


Entanglement/Squeezing manipulation

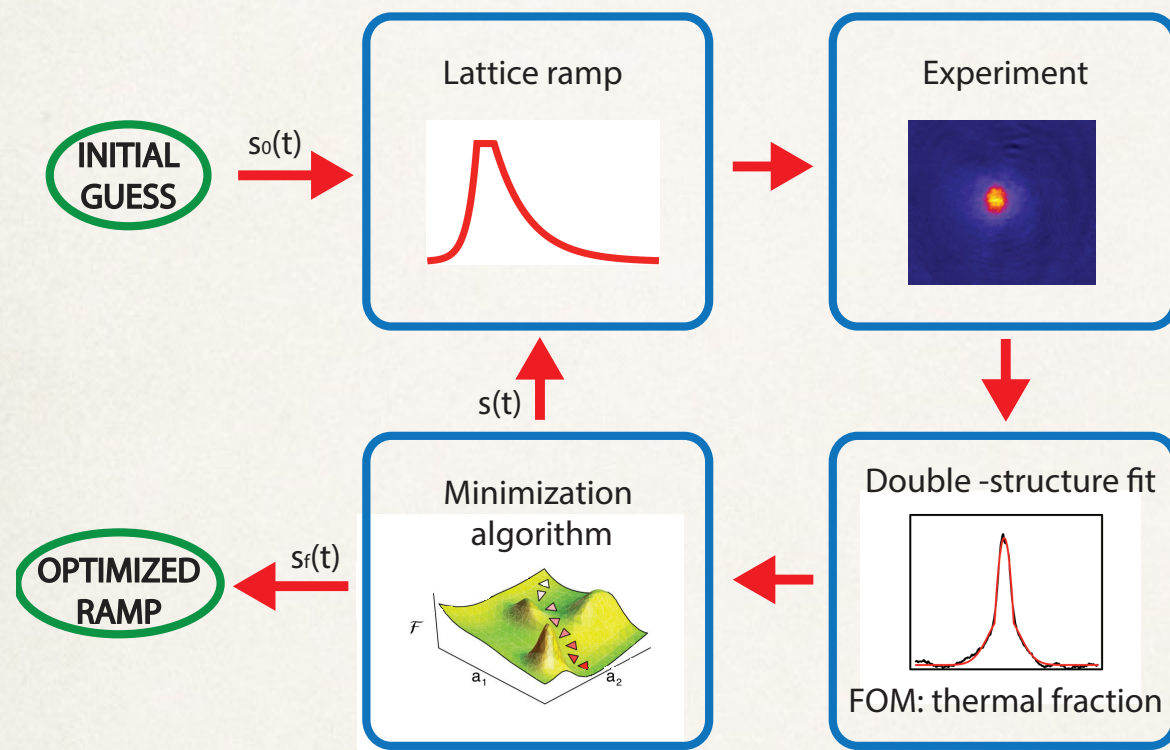


Optimal experimental protocols

Atom chip interferometer at QSL



Closed loop optimization

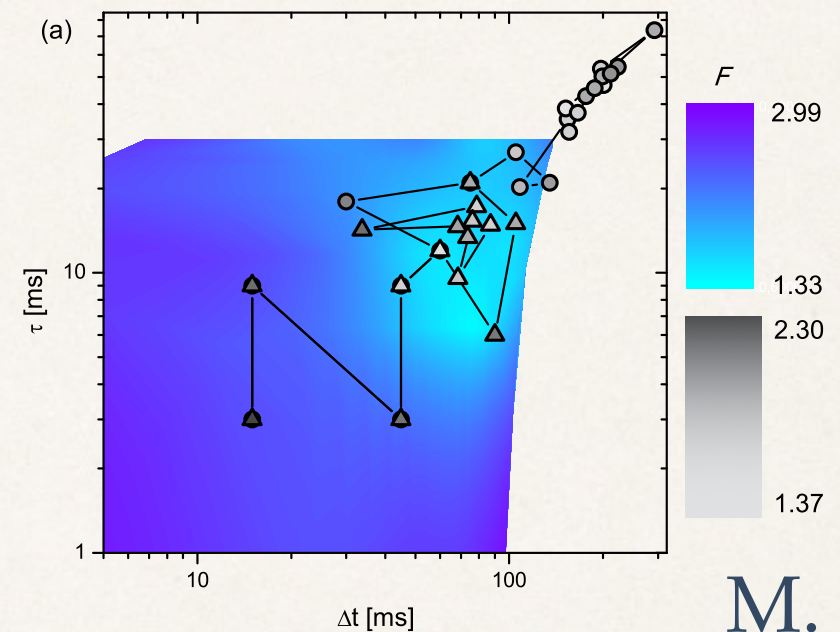


3D-1D crossover and QPT

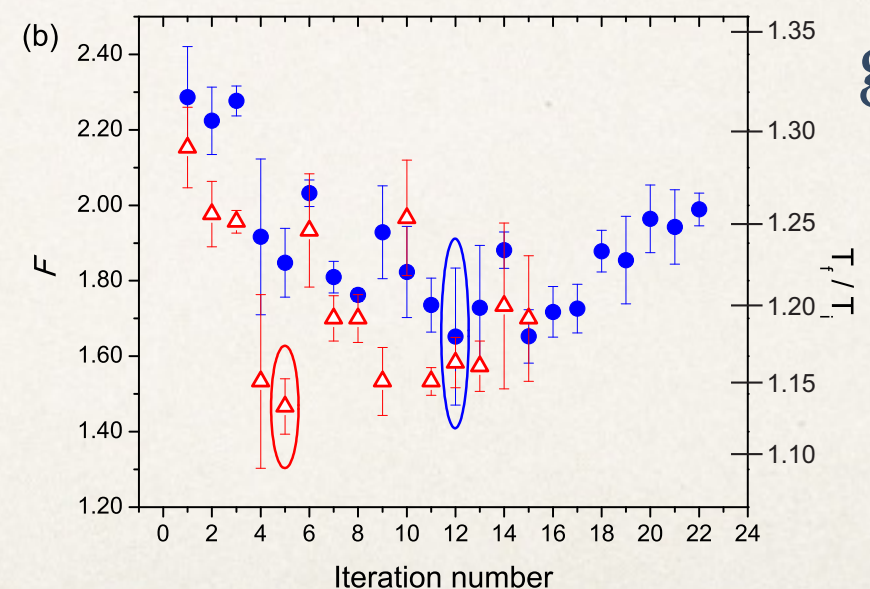
$$T_{opt} \sim T_{ad}/3$$

$$FOM_{opt} \sim 0.9 FOM_{ad}$$

S. Rosi, et. al. PRA 2013



M. Inguscio
group



Run	Δt_{opt}	τ_{opt}	F_{uncorr}	Best F_{opt}
1	154 ms	35 ms	2.30 ± 0.03	1.73 ± 0.02
2	45 ms	9 ms	2.30 ± 0.03	1.40 ± 0.06

Conclusions

- ❖ CRAB optimization can be applied successfully to MBQS dynamics opening new perspectives.
- ❖ What can be simulated can be controlled
- ❖ Optimal trajectories are robust with respect to noise and perturbations.
- ❖ Non-integrable MBQS are exponentially complex to be optimized (in general) but particular interesting protocols can be efficiently simulated and thus controlled.



Cartagena (Spain) 25-30/5/2015

New Trends in Complex Quantum systems dynamics - 2nd Edition

$c(t)$

Thank you for your attention!

Tommaso Calarco
Tommaso Caneva
Matthias Mueller
Thomas Pichler
Enrique Rico
Martin Plenio
Fedor Jelezko
Boris Nayadov



Massimo Inguscio
Leonardo Fallani
Alain Bernard
Chiara Fort
Nicole Fabbri
Filippo Caruso



Immanuel Bloch
Marc Cheneau
Sebastian Hild



Misha Lukin



Tilman Pfau
Robert Löwe



Peter Zoller
Marcello Dalmonte



S. Lloyd



Rosario Fazio



Jörg Schmiedmayer
Thorsten Schumm
Sandrine van Frank
Wolfgang Rohringer



Funds:

SFB / TRR21 Co.Co.Mat.



IP-SIQS
IP-DIADEMS
STREP-DIAMANT
STREP-PICC
STREP-MALICIA

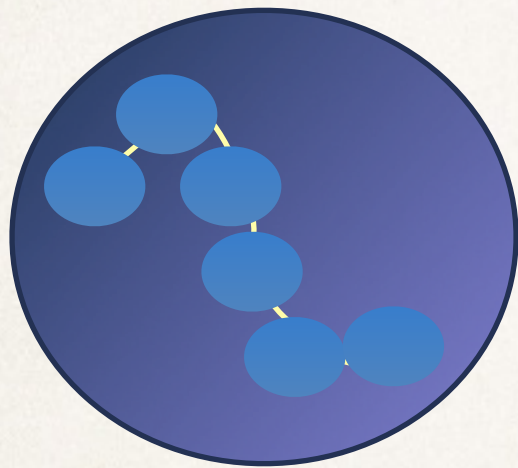


Numerics:

BW-Grid

www.dmrg.it

Upper-bound



$$n_\varepsilon = L/\varepsilon \leq T v_{max}/\varepsilon = \text{Poly}(D_{\mathcal{W}+}) v_{max}/\varepsilon$$

Bounded energy

T-poly

$$b_S^+ = \frac{\text{Poly}(D_{\mathcal{W}+}) v_{max}}{\varepsilon} \log \varepsilon^{-D_{\mathcal{W}+}}$$

$$b_\gamma \leq b_S^+$$

$$\text{Poly}'(D_{\mathcal{W}+}) v_{max}/\varepsilon \geq n_s$$

In conclusion:

$$n_s = \text{Poly}(D_{\mathcal{W}+})$$

Smoothed complexity

Analysis of algorithmic complexity **under perturbed input**

Machine learning, Numerical analysis, Discrete mathematics...

Linear programming:

$$\begin{aligned} &\max c^T x \\ &\text{subject to} \\ &Ax \leq b \end{aligned}$$

Simplex algorithm has polynomial smoothed complexity!